

The description and the verification of EAM-potentials at various electron temperatures

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The absorption of laser pulse initially leads to the excitation of electron subsystem (ES). In this case initial state of the system is two-temperature (2T) state, and the electron temperature (T_e) may be several orders higher than the ion temperature (T_i). For gold the change of interionic forces plays the key role in dynamics of ions at the initial 2T-stage. One of the ways to describe such process is to use electron-temperature-dependent (ETD) interionic potential in a simulation. Such potential for gold was developed in works [1-3].

For creation of ETD-potential for gold, three EAM-potentials (potentials in the Embedded Atom Method form) at different electron temperatures were created. Construction of the EAM-potentials was performed with force-matching method [4,5]. This technique was used as implemented in the Potfit code [5]. The method provides a way to construct physically justified interparticle potentials without referring to experimental data. The idea is to adjust the interparticle potential to optimally reproduce per-atom forces computed at the *ab initio* level (e.g. with Density Functional Theory) for a fine-tuned set of small reference structures. The reference data were calculated using the VASP code [6]. In our case, the calculations of the reference data were performed at three different T_e (0.1, 3 and 6 eV) that are set as a parameter of the Fermi-Dirac distribution for partial occupancies of electron bands. All the three developed EAM-potentials reproduce forces calculated by the VASP code to within 15–20%, which is quite good accuracy for the given method. The EAM-potential has three independent functions: pair interaction $\varphi(r)$, effective electron density $\rho(r)$ and embedding energy term Φ . Fig. 1 shows the functions of the created EAM-potentials.

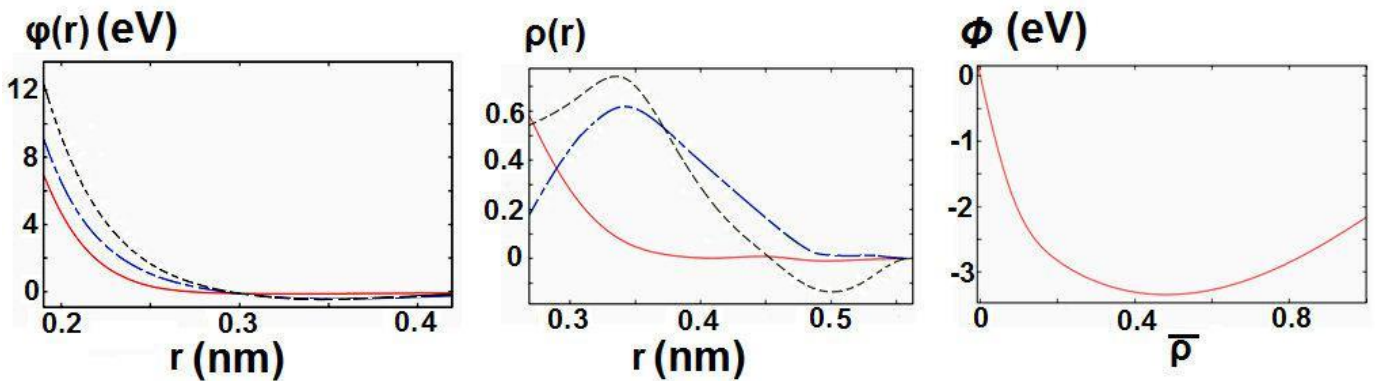


Fig 1. The created EAM-potential functions: $\varphi(r)$ and $\rho(r)$ for gold at various electron temperatures $T_e = 0.1$ eV (solid curves), 3 eV (dash-dot curves), 6 eV (dashed curves). The “embedding” function $\Phi(\rho)$ is shown for $T_e = 0.1$ eV.

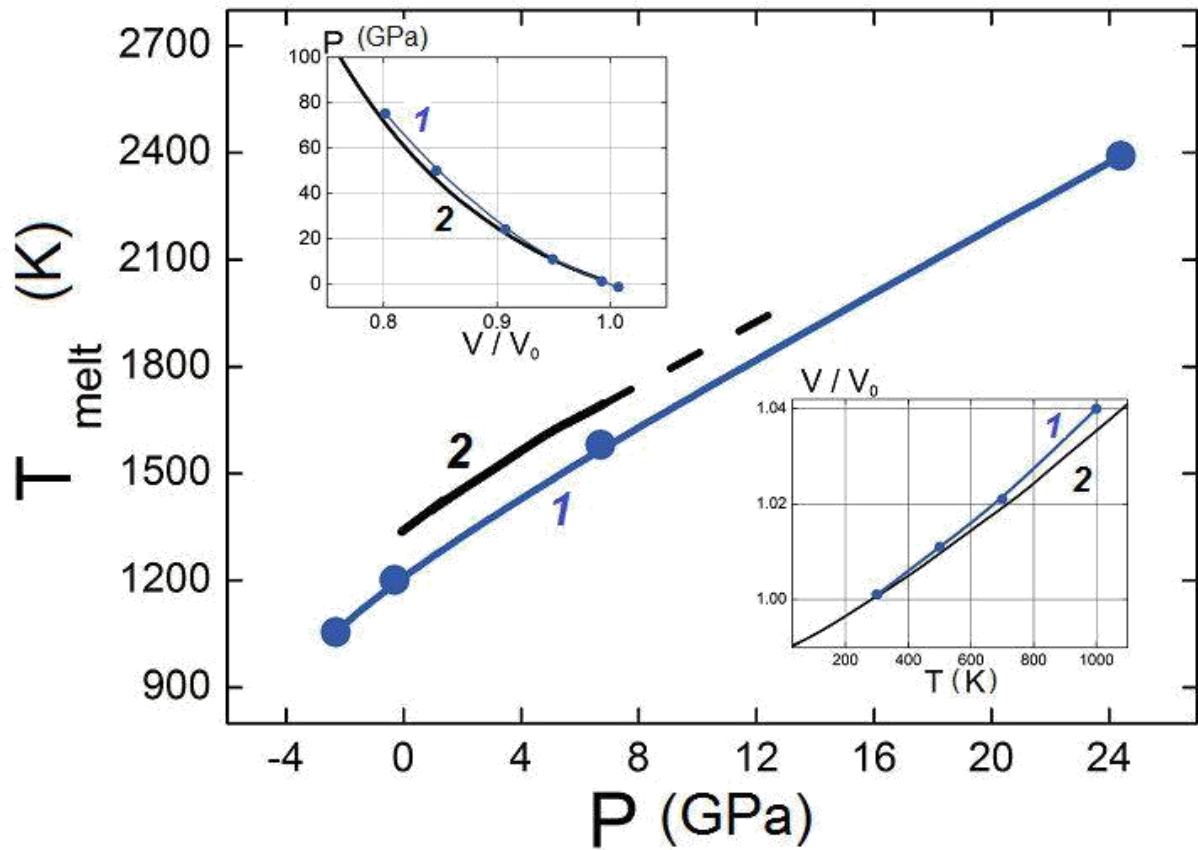


Fig 2. Plots of the melting temperature T_{melt} versus pressure P , room-temperature P versus V isotherm (top inset), and thermal expansion V/V_0 (bottom inset) of gold: (1) simulation using the proposed EAM potential ($T_e = 0.1$ eV); (2) experimental data [7–9] (V_0 is equilibrium volume)

Fig. 2 and Table show the results of verification of the EAM-potential at $T_e = 0.1$ eV (*cold part* of ETD-potential). All created EAM-potentials are given in **Au_5T.eam.alloy** file (eam/alloy format for LAMMPS). The indexes **Au**, **Au_3.0**, **Au_6.0** indicate three initially developed EAM-potential: **Au** for EAM-potential at $T_e = 0.1$ eV (*cold part*), **Au_3.0** **Au** for EAM-potential at $T_e = 3.0$ eV, **Au_6.0** **Au** for potential at $T_e = 6.0$ eV. The indexes **Au_1.5** and **Au_4.5** indicate EAM-potential which were created by interpolation. As the potentials are given in format for alloy, it is necessary to create the cross-pair interactions for various types. All cross-pair interactions are given in **Au_5T.eam.alloy** file as the average pair interactions. For instance the cross pair potential $\varphi(r)$ for "1.5" and "6.0" types is $0.5 \cdot (\varphi_{1.5} + \varphi_{6.0})$.

Table. The simulated properties of gold with using of EAM-potential at $T_e = 0.1$ eV. The experimental data [7–9] are shown in brackets.

V_0 (cm ³ mol ⁻¹)	E_c (eV)	C_{11} (GPa)	C_{12} (GPa)	T_{melt} (K)
10.23 (10.22)	4.1 (3.8)	230 (202)	180 (170)	1210 (1338)

The properties at high electron temperatures

The ETD-potential may be created using quadratic polynomial interpolation [3]:

$$\varphi(r) = \varphi_0 + \varphi_1 T_e + \varphi_2 T_e^2$$

and

$$\rho(r) = \rho_0 + \rho_1 T_e + \rho_2 T_e^2.$$

The ETD-potential coincides with the EAM-potentials at three reference temperatures. At increasing of T_e the ETD-potential may describe the features of 2T-system. In particular, one of the phenomena that can be reproduced by the ETD-potential is an increase of the ion melting temperature T_{melt} at heating of the ES [10]. Fig. 3 shows the results of the simulation.

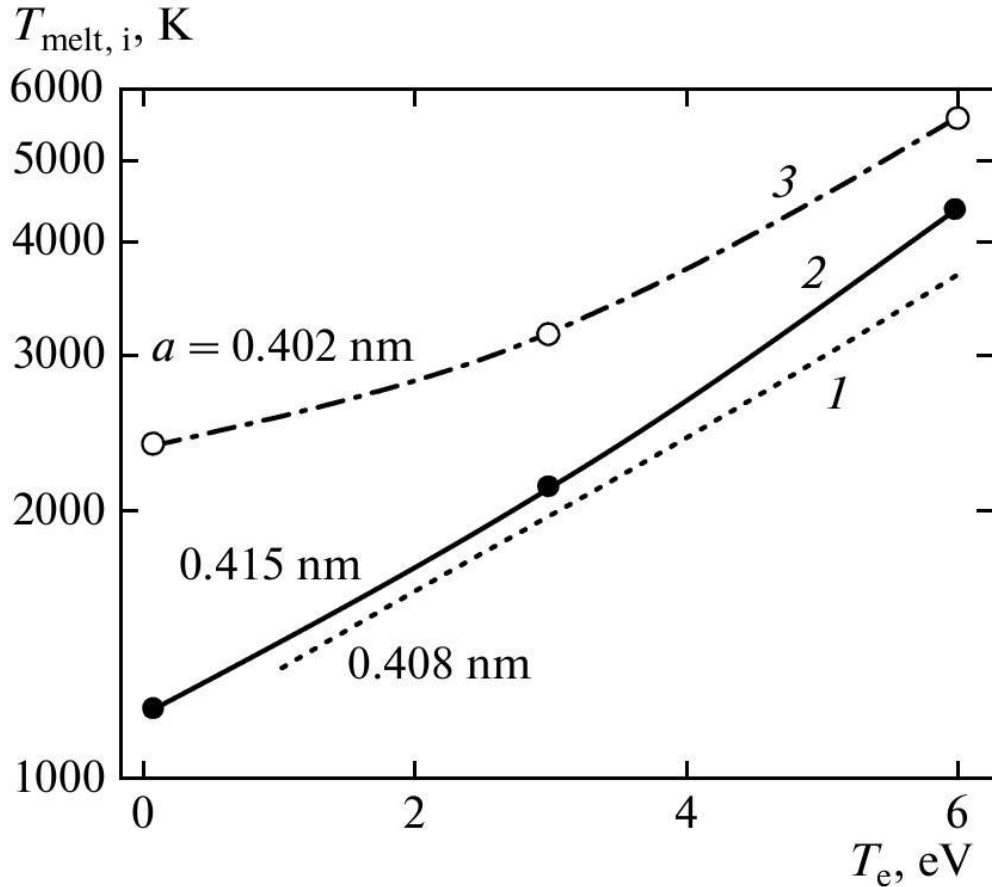


Fig 3. Plots of ion melting temperature T_{melt} versus T_e for gold at various densities (lattice parameters a): (1) data from [10]; (2, 3) results of two-phase atomistic simulations in this study for EAM-potentials at $T_e = 0.1, 3,$ and 6 eV.

Another phenomenon that is reproduced by the ETD-potential is increase of the pressure P_e^{loc} of localized electron energy with increase of T_e . In fact, as T_e increases from 0 to 6 eV the pressure in the system described by the ETD-potential increases from 0 to 105 GPa. Together with the value of P_e^{deloc} the total pressure in gold at $T_e = 6$ eV reaches 200 GPa like in work [10]. Figure 4 shows plots of the total pressure P (*ab initio* calculations using the VASP code) and P_e^{loc} (MD-calculation with the ETD-potential) versus T_e . The difference between the two pressures determines the dependence $P_e^{deloc}(T_e)$.

Some verification of T_e -part of the potentials (about electron pressure) was given in works [11, 12]

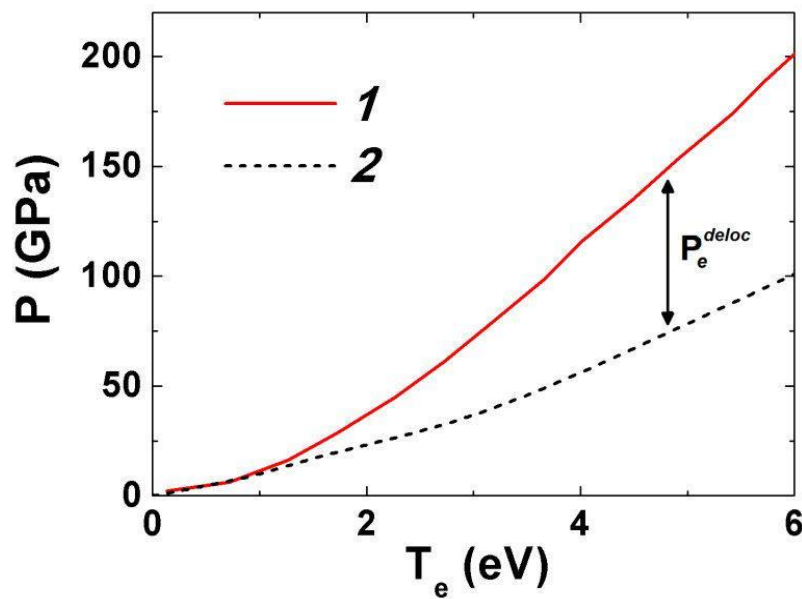


Fig 4. The dependence of pressure on T_e : 1 – total pressure P (*ab initio* calculations with using the VASP code); 2 – the P_e^{loc} reconstructed with the virial theorem (molecular dynamics calculation with the ETD potential). The difference between the two pressures determines the dependence $P_e^{deloc}(T_e)$.

References

- [1] G. Norman, S. Starikov, V. Stegailov, V. Fortov, et al., J. Appl. Phys. **112**, 013104 (2012)
- [2] G.E. Norman, S.V. Starikov, V.V. Stegailov, Journal of Experimental and Theoretical Physics, **114**, No. 5, pp. 792–800 (2012)
- [3] S. Starikov, A. Faenov, T. Pikuz, I. Skobelev, et al., Appl. Phys. B. **116**, P. 1005-1016 (2014)
- [4] F. Ercolessi, and J.B. Adams. Europhysics Letters, **26**, 583-588 (1994).
- [5] P. Brommer, and F. Gahler. Philosophical Magazine, **86**, 753-758 (2006).
- [6] G. Kresse and J. Furthmüller, Phys. Rev. B **54**, 11169–11186 (1996).
- [7] J.R. Neighbours, G.A. Alers, Phys. Rev. **111**, P.707 (1958)

[8] P.M. Bell, J. Xu, H.K. Mao, in Shock Waves in Condense Matter, edited by Y.M. Gupta (Plenum, NewYork, 1986).

[9] I.-K. Suh, H. Ohta, Y. Waseda, J. Mater. Sci. **23**, P.757 (1988)

[10] V. Recoules et al. Phys. Rev B, **96**, 05550 (2006)

[11] V. Stegailov and P. Zhilyaev, Pressure in electronically excited warm dense metals, Contrib. Plasma Phys. **55**, 164 – 171 (2015)

[12] V. Stegailov and P. Zhilyaev, Warm dense gold: effective ion–ion interaction and ionization, MOLECULAR PHYSICS (2015)